

## context

- given a mapping  $m$  from  $(\hat{S}, \hat{g}_{\hat{a}\hat{b}})$  to  $(S, g_{ab})$
- for parameterization of a disk like surface  $S$ ,  $\hat{S}$  will be a disk-like portion of the plane.
- in practice  $S$  will be a triangle mesh
- we will map each mesh triangle to one triangle in the plane ‘(One Piece linear) ‘O.P.L.’”
- we will think of the positions of the vertices in the plane as the eventual unknowns.

## pullback

- we have the differential tangent map  $dm_a^b$
- we define the pull back of  $g$  as

$$g_{\hat{a}\hat{b}} \equiv dm_b^d g_{cd} dm_a^c$$

- this allows us effectively apply  $g$  to tangent vectors of  $\hat{S}$ .

## isometry

- we call a mapping  $m$  from  $(\hat{S}, \hat{g}_{\hat{a}\hat{b}})$  to  $(S, g_{ab})$  isometric if

$$\hat{g}_{\hat{a}\hat{b}} = g_{ab}$$

- this means that the dot product is preserved under the mapping.
- so are lengths, and so are all “distances” between points
- most pairs of surfaces do not have an isometric mapping between them

## conformal

- we call a mapping  $m$  from  $(\hat{S}, \hat{g}_{\hat{a}\hat{b}})$  to  $(S, g_{ab})$  conformal if

$$\lambda^2 \hat{g}_{\hat{a}\hat{b}} = g_{ab}$$

- for some spatially varying  $\lambda$
- this means that at a point, all dot products are just scaled
- this implies that angles are preserved
- this constraint places  $3 - 1 = 2$  constraints on the mapping at each point
  - symmetric 2 by 2 matrix has 3dof, scale has 1
- deep fact: there is a conformal map between any two disc-like surfaces
  - unique up to 3 dofs.
  - so you cannot fix the mapping at the boundary.
  - if we require the map to be O.P.L. we will most certainly fail

## benefits of conformal

- preserves “shape”, but areas are not controlled
- if sampling rate can be adaptive but isotropic (quad tree), then can overcome the area distortion.
- preserves aspect ratios of triangles
- lots of continuous theory

## area preserving

- we call the map area preserving if

$$dA_g = dA_{\hat{g}}$$

- this imposes one constraint on the mapping at each point
- fact: there are loads of area preserving maps between pairs of disks.
- e.g. you can fix the boundary mapping
- you can add other criteria to optimize as well
- if you require the map to be O.P.L. you may or may not fail.
- not pursued that much in the literature
- this is also useful for cartograms.

## stretch

- define the (squared) stretch of a vector  $v^{\hat{a}}$  as

$$\frac{v^{\hat{a}} g_{\hat{a}\hat{b}} v^{\hat{b}}}{v^{\hat{a}} \hat{g}_{\hat{a}\hat{b}} v^{\hat{b}}}$$

- it is the size of the unitized vector under  $m$
- we wish to understand the stretch for various directions
- we want to figure out the directions of max and min stretch
- this will lead to a notion of non-conformal energy to minimize

## dirchilet tensor

- define the “dirchilet” tensor as

$$\begin{aligned} p_{\hat{e}}^{\hat{a}} &\equiv \hat{g}^{\hat{a}\hat{b}} dm_b^c g_{cd} dm_{\hat{e}}^d \\ &= \hat{g}^{\hat{a}\hat{b}} g_{\hat{b}\hat{e}} \end{aligned}$$

- $p_{\hat{e}}^{\hat{a}}$  is self adjoint wrt  $\hat{g}_{\hat{a}\hat{b}}$  since  $p_{\hat{e}}^{\hat{a}} \hat{g}_{\hat{a}\hat{b}} = g_{\hat{b}\hat{e}}$  is symmetric
- so by the spectral thm,  $p$  has 2 eigenvectors that are ortho wrt  $\hat{g}$ .
- these are the max and min of

$$\frac{v^{\hat{a}} g_{\hat{a}\hat{b}} v^{\hat{b}}}{v^{\hat{a}} \hat{g}_{\hat{a}\hat{b}} v^{\hat{b}}} = \frac{\hat{g}_{\hat{a}\hat{c}} p_{\hat{b}}^{\hat{c}} v^{\hat{a}} v^{\hat{b}}}{\hat{g}_{\hat{a}\hat{b}} v^{\hat{a}} v^{\hat{b}}}$$

## properties of p

- call the two orthogonal (wrt  $\hat{g}$ ) eigenvectors  $n_i^{\hat{e}}$  with positive eigenvalues  $\gamma_1^2 \geq \gamma_2^2$
- stretch is maximized for vectors in direction  $n_1^{\hat{e}}$  and minimized for  $n_2^{\hat{e}}$
- easy fact: we have  $\frac{dA_g}{dA_{\hat{g}}} = \gamma_1 \gamma_2 = \sqrt{DET(P)}$ 
  - where  $P$  is the matrix expression of  $p$  using any set of coordinates

## non-conformality energy

- for a conformal map, we must have  $\gamma_1 = \gamma_2$

- so one way to measure non-conformality is

$$\begin{aligned} & \int_{\hat{S}} dA_{\hat{g}} (\gamma_1 - \gamma_2)^2 = \\ & \int_{\hat{S}} dA_{\hat{g}} (\gamma_1^2 + \gamma_2^2) - 2 \int_{\hat{S}} dA_{\hat{g}} (\gamma_1 \gamma_2) = \\ & \int_{\hat{S}} dA_{\hat{g}} (\gamma_1^2 + \gamma_2^2) - 2 \int_{\hat{S}} dA_g = \\ & \int_{\hat{S}} dA_{\hat{g}} (\gamma_1^2 + \gamma_2^2) - 2A(S) \end{aligned}$$

- if we are optimizing this energy over maps where  $S$  is fixed, then so is its area, and so we may as well just optimize the “dirichilet” energy:

$$P = \int_{\hat{S}} dA_{\hat{g}} (\gamma_1^2 + \gamma_2^2)$$

- dirichilet energy  $\equiv$  non-conformality + 2 area of range

### variational

- we can look at some space of mappings and minimize the dirichilet energy over this space
- if we lock the down the mapping along the boundary, we cannot get a conformal map.
- the minimizer is called “harmonic”

### dirichilet energy

- we can compute the  $\gamma_i$  by picking any coordinates, representing  $p$  as a matrix, and diagonalizing  $P = E^{-1}DE$
- the quantity  $(\gamma_1^2 + \gamma_2^2)$  is trace(D).
- but trace(D)=trace(P)=  $p_a^{\hat{a}}$ .
- so we can also write the dirichilet energy as

$$\int_{\hat{S}} dA_{\hat{g}} p_a^{\hat{a}} = \int_{\hat{S}} dA_{\hat{g}} \hat{g}^{\hat{a}\hat{b}} dm_b^c g_{cd} dm_a^d$$

### conformal invariant

- we are looking at the dirichilet energy of a mapping from  $\hat{S}$  to  $S$
- if I also have a conformal mapping from  $\hat{\hat{S}}$  to  $\hat{S}$ .
- this gives me a composed map from  $\hat{\hat{S}}$  to  $S$ .
- fact: the energy of the composed map is the same as the energy of the original map
- related to the following: swap the metric  $\hat{\hat{g}} = \lambda^2 \hat{g}_{\hat{a}\hat{b}}$
- so dirichilet integrand (which has a factor of  $\hat{\hat{g}}^{\hat{a}\hat{b}}$ ) scales with  $1/\lambda^2$ , but  $dA_{\hat{\hat{g}}^{\hat{a}\hat{b}}}$  scales with  $\lambda^2$
- relevance: if  $(\hat{\hat{S}})$  is allowed to scale, the optimization will not im/explode it.

### MIPS (corrected?)

- if the mapping is P.L., then  $dm$  is constant per triangle so dirichilet energy of the mapping is

$$\begin{aligned} P &= \int_{\hat{S}} dA_{\hat{g}} (\gamma_1^2 + \gamma_2^2) \\ &= \sum_{t \in T} A(t) (\gamma_1^2(t) + \gamma_2^2(t)) \end{aligned}$$

- where  $A(\hat{t})$  is the area of the triangle in the plane.
- this will end up being some non-quadratic expression in the  $(u_i)$  positions of the vertices.
- non-linear numerical optimization
- for continuous case, any conformal map to any region of the plane will have Dirichlet energy  $P = 2A(S)$ .
- for OPL, we will not have any conformal maps, so we will have some non-conformal energy.
  - so some boundary shapes in the plane will lead to less energy than others
- but due to conformal invariance, it will not try to implode so one can just leave the boundaries free
- cost goes infinite at degenerate area triangle, so will avoid flips (i think).
- i call this MIPS
  - MIPS as originally published is a bit different, and did not take into account the triangles areas, and so is not intrinsic.

## Eck

- we also have an inverse mapping  $u := m^{-1}$  from  $S$  to  $\hat{S}$
- one can compute the Dirichlet energy of that inverse mapping
- the extrema of the stretch of the inverse mapping are  $\lambda_i = 1/\gamma_i$
- we will still have  $\hat{S}$  be the plane.
- so the Dirichlet energy of the inverse mapping can be computed as

$$\int_S dA_g (\lambda_1^2 + \lambda_2^2) = \sum_{t \in T} A(t) (\lambda_1^2 + \lambda_2^2)$$

- where  $A(t)$  is 3d area of the triangle.

## Eck properties

- -: Dirichlet = non-conformal + area of range
  - so Dirichlet of inverse mapping will want to implode the parameter domain
- so fix the boundary
- could allow to slide along boundary, but this would be complicated to constrain
- so just fix the boundary mapping
- we won't get a conformal map, but the optimal is called a "harmonic map".
- in regions that are very non-conformal, Eck will try to shrink, so there are some quality issues
- this is called Eck.

## Eck as Poisson problem

- $\hat{S}$  is the plane, so let's use planar coordinates  $u_1, u_2$ .
- then the energy is

$$\int_S dA_g g^{ab} du_b^{\hat{c}} \hat{g}_{\hat{c}\hat{d}} du_a^{\hat{d}} = \sum_{\hat{i}} \int_S dA_g g^{ab} du_b^{\hat{i}} du_a^{\hat{i}}$$

- this is exactly the same energy from the sum of 2 Poisson problems, with a zero right hand side!!

## PL case

- since  $S$  is a mesh, we can write this as

$$\begin{aligned} \sum_i ((u^i)^t d^t M_2 du^i) &= \sum_i 1/2 \sum_t \sum_{he_{vw}} (\cot \alpha_{vw}) (u_v^i - u_w^i)^2 \\ &= \sum_i 1/2 \sum_{e_{vw}} (\cot \alpha_{wv} + \cot \alpha_{vw}) (u_v^i - u_w^i)^2 \end{aligned}$$

- +: the energy is quadratic in the vertex position
- can be written as spring energy with symmetric “cotangent weights”
  - lots of theory for this
  - can be negative weights (so possible flips).

## levy-Desbrun

- go back to the original non-conformal energy of the inverse mapping

$$\int_S dA_g (\lambda_1 - \lambda_2)^2$$

- this likes conformality. if not conformal, will try to shrink, but not at the expense of conformality.
- in smooth case, any conformal map will have zero energy
- for O.P.L. there is no conformal map, and so minimum solution is a full collapse.
  - unless it was flat to begin with
- must add at least 2 fixed vertex constraints. (choice changes the result).
- this will not want to collapse
- this is still a linear system (see below)!
  - so is strictly better than Dirichlet.
  - for O.P.L., we still get cotangent springs at interior vertices and linear but non separable equations at the boundary vertices
- linear system, may have flips due to negative weights, or due to the non-convexity of the boundary.

## as similar as possible

- now, to obtain another interpretation for levy....
- lets set up a new problem
- for each mapping  $u$ , and spatially varying similarity transform  $C$ , define an energy as

$$\int_S dA_g (du_b^{\hat{a}} - C_b^{\hat{a}})(du_d^{\hat{c}} - C_d^{\hat{c}})g^{bd}\hat{g}_{\hat{a}\hat{c}}$$

- since  $\hat{S}$  is the plane, we get

$$\sum_{\hat{k}} \int_S dA_g (du_b^{\hat{k}} - C_b^{\hat{k}})(du_d^{\hat{k}} - C_d^{\hat{k}})g^{bd}$$

- minimize over  $u$  and  $C$ .
- if  $u$  is conformal mapping, there will be a  $C$  that gives us zero energy
- this is a sum of two poisson energies, where the  $C^k$  is the right hand side.

## rewrite

- so for ASAP, on a triangle mesh, we get

$$\begin{aligned} & \sum_{\hat{k}} 1/2 \sum_t \sum_{he_{vw}} (\cot \alpha_{vw}) ((u_v^{\hat{k}} - u_w^{\hat{k}}) - \sum_l C_l^{\hat{k}}(t) [x_v^l - x_w^l])^2 \\ & = 1/2 \sum_t \sum_{he_{vw}} (\cot \alpha_{vw}) \|((u_v - u_w) - C(t)[x_v - x_w])\|^2 \end{aligned}$$

- in the last expression,  $C(t)$  is a 2-by-2 matrix

### PL ASAP

- for P.L. mapping, and one  $C$  per triangle
  - parameterize as  $[[a, b], [-b, a]]$
- we get a linear system in our variables!

### interpretation of ASAP

- $C$  is just a helper variable. we want to understand this as an optimization over the mapping  $u$ .
- define  $E(u) = \min_C E(u, C)$
- then  $\min_{u, C} E(u, C) = \min_u E(u)$
- if we understand  $E(u)$ , then we understand ASAP.
- define  $C^*(u) = \operatorname{argmin}_C E(u, C)$ ,
- then  $E(u) = E(u, C^*(u))$
- so if we understand  $C^*(u)$ , we understand ASAP.

### procrustus

$$\min_C \int_S dA_g (du_b^{\hat{a}} - C_b^{\hat{a}})(du_d^{\hat{c}} - C_d^{\hat{c}}) g^{bd} \hat{g}_{\hat{a}\hat{c}}$$

- fixed  $du$ , and  $C$  must be a similarity transform
- can just look at each point (or triangle) at a time
- and fix parameters at this point so that both sides are orthonormal and get a matrix problem

$$\min_C |du_j^{\hat{k}} - C_j^{\hat{k}}|_F^2$$

- note that in these coords,  $P^{-1} = du^t * du$ .
- by (signed) SVD:

$$du_j^{\hat{k}} = U \begin{bmatrix} \lambda_1 & 0 \\ 0 & \lambda_2 \end{bmatrix} V^t$$

- one  $\lambda$  is negative if  $du$  is a flip
- by Procrustus:

$$C^* = \frac{1}{2} U \begin{bmatrix} \lambda_1 + \lambda_2 & 0 \\ 0 & \lambda_1 + \lambda_2 \end{bmatrix} V^t$$

- if you plug this  $C^*$  you get

$$E(u) = \int_S dA_g (\lambda_1 - \lambda_2)^2$$

– use the permuted trace rule (see paper)

- thus we are back to Levy-Desbrun.

## ARAP

- if we replace the similarity  $C$ , with a rotation  $R$ , we will get “as rigid as possible” parameterization
- by Procrustus:

$$R^* = UV^t$$

- plug this in

$$\int_S dA_g (\lambda_1 - 1)^2 + (\lambda_2 - 1)^2$$

## ARAP computation

$R$  is subject to non linear constraints, so this will be a non linear problem

in PL case, if we fix  $R$ , we get a poisson system to find best  $u$ . and for fixed  $u$  can find best  $R$  per triangle using procrustus so can iterate back and forth.

## green energy

- measures internal stresses
- tries to find an isometry
- generally it cannot succeed

$$\begin{aligned} & \int_S dA_g |\hat{g}_{ab} - g_{ab}|^2 \\ &= \int_S dA_g (\hat{g}_{ab} - g_{ab})(\hat{g}_{cd} - g_{cd})g^{ac}g^{bd} \\ &= \int_S dA_g (1 - \lambda_1^2)^2 + (1 - \lambda_2^2)^2 \end{aligned}$$

- non linear
- doesn't mind flips